

## REVIEW OF CHARM DALITZ PLOT ANALYSES

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For references given here in the form SMITH 05, see the references at the end of the  $D^+$ ,  $D^0$ , and  $D_s^+$  Listings.

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The formalism of Dalitz-Plot analysis is reviewed in the preceeding note. Table 1 lists reported analyses of  $D$  mesons. In the following, we discuss a number of subjects of current interest: (1)  $D^0 \rightarrow K_S^0 \pi^+ \pi^-$ ; (2)  $D \rightarrow \pi \pi \pi$ : a  $\sigma(500)$  or  $f_0(600)$ ; (3)  $D^+ \rightarrow K^- \pi^+ \pi^+$ : a  $\kappa(800)$ ? (4) the  $f_0(980)$ ,  $f_0(1370)$  and  $f_0(1500)$ ; (5) doubly Cabibbo-suppressed decays; and (6)  $CP$  violation.

$D^0 \rightarrow K_S^0 \pi^+ \pi^-$ : Several experiments have analyzed  $D^0 \rightarrow K_S^0 \pi^+ \pi^-$  decay (see Table 1). The most precise results are from CLEO (BABAR and Belle, discussed below, have not yet evaluated systematic uncertainties). The CLEO analysis included ten resonances:  $K_S^0 \rho^0$ ,  $K_S^0 \omega$ ,  $K_S^0 f_0(980)$ ,  $K_S^0 f_2(1270)$ ,  $K_S^0 f_0(1370)$ ,  $K^*(892)^- \pi^+$ ,  $K_0^*(1430)^- \pi^+$ ,  $K_2^*(1430)^- \pi^+$ ,  $K^*(1680)^- \pi^+$ , and the doubly Cabibbo-suppressed mode  $K^*(892)^+ \pi^-$ . CLEO found a much smaller nonresonant contribution than did the earliest experiments.

The source of the nonresonant component found in the early experiments has been attributed to the broad scalar resonances, the  $K_0^*(1430)^-$  and  $f_0(1370)$ , found in the later, larger data samples. The observation of a small but significant nonresonant component in the largest data samples suggests the presence of additional broad scalar resonances, the  $\kappa(800)$  and  $\sigma(500)$ . The CLEO analysis could accommodate the  $\sigma(500)$  in lieu of the nonresonant component, but found no evidence for the  $\kappa(800)$ .

The ten quasi-two-body intermediate states in the CLEO analysis include both  $CP$ -even and  $CP$ -odd eigenstates and one doubly Cabibbo-suppressed channel. A time-dependent analysis of the Dalitz plot allows simultaneous determination of the strong transition amplitudes and phases and the mixing parameters  $x$  and  $y$  without phase or sign ambiguity. Using  $9 \text{ fb}^{-1}$ , CLEO obtained  $(-4.5 < x < 9.3)\%$  and  $(-6.4 < y < 3.6)\%$  [1].

The decay  $D^0 \rightarrow K_S^0 \pi^+ \pi^-$ , important for the study of the CKM angle  $\gamma/\phi_3$  [6], is under study by Belle [10,11] and BABAR [12,13]. The CLEO model does not provide a good description of the higher-statistics BABAR and Belle data samples. An improved description is obtained in two ways: First, by adding more Breit-Wigner resonances, including two  $\pi\pi$  resonances with arbitrary mass and width, denoted as  $\sigma_1$

**Table 1:** Reported Dalitz plot analyses.

Decay	Experiment(s)
$D^0 \rightarrow K_S^0 \pi^+ \pi^-$	Mark II <sup>a</sup> , Mark III <sup>b</sup> , E691 <sup>c</sup> , E687 <sup>d,e</sup> , ARGUS <sup>f</sup> , CLEO <sup>g</sup> , Belle [10,11], BABAR [12,13]
$D^0 \rightarrow K^- \pi^+ \pi^0$	Mark III <sup>b</sup> , E687 <sup>e</sup> , E691 <sup>c</sup> , CLEO <sup>h</sup>
$D^0 \rightarrow \bar{K}^0 K^+ \pi^-$	BABAR [14]
$D^0 \rightarrow K^0 K^- \pi^+$	BABAR [14]
$D^0 \rightarrow K_S^0 \eta \pi^0$	CLEO <sup>i</sup>
$D^0 \rightarrow \pi^+ \pi^- \pi^0$	CLEO <sup>j</sup>
$D^0 \rightarrow K_S^0 K^+ K^-$	BABAR <sup>k</sup>
$D^0 \rightarrow K^- K^+ K^- \pi^+$	FOCUS <sup>l</sup>
$D^0 \rightarrow K^- K^+ \pi^- \pi^+$	FOCUS <sup>m</sup>
$D^+ \rightarrow K^- \pi^+ \pi^+$	Mark III <sup>b</sup> , E687 <sup>e</sup> , E691 <sup>c</sup> , E791 <sup>n</sup>
$D^+ \rightarrow \bar{K}^0 \pi^+ \pi^0$	Mark III <sup>b</sup>
$D^+ \rightarrow \pi^+ \pi^+ \pi^-$	E687 <sup>o</sup> , E791 <sup>p</sup> , FOCUS [5] <sup>q</sup>
$D^+ \rightarrow K^+ K^- \pi^+$	FOCUS [15], E687 <sup>r</sup> , BABAR <sup>s</sup>
$D^+ \rightarrow K^+ \pi^+ \pi^-$	E791 <sup>t</sup> , FOCUS <sup>u</sup>
$D_s^+ \rightarrow K^+ K^- \pi^+$	E687 <sup>r</sup> , FOCUS [15]
$D_s^+ \rightarrow \pi^+ \pi^+ \pi^-$	E687 <sup>o</sup> , E791 <sup>v</sup> , FOCUS [5]
$D_s^+ \rightarrow K^+ \pi^+ \pi^-$	FOCUS <sup>u</sup>

See the end of the  $D^+$ ,  $D^0$  and  $D_s^+$  Listings for these references:

<sup>a</sup>SCHINDLER 81, <sup>b</sup>ADLER 87, <sup>c</sup>ANJOS 93, <sup>d</sup>FRABETTI 92B, <sup>e</sup>FRABETTI 94G, <sup>f</sup>ALBRECHT 93D, <sup>g</sup>MURAMATSU 02, <sup>h</sup>KOPP 01, <sup>i</sup>RUBIN 04, <sup>j</sup>CRONIN-HENNESSY 05, <sup>k</sup>AUBERT 05B, <sup>l</sup>LINK 03G, <sup>m</sup>LINK 05C, <sup>n</sup>AITALA 02, <sup>o</sup>FRABETTI 97D, <sup>p</sup>AITALA 01B, <sup>q</sup>LINK 04, <sup>r</sup>FRABETTI 95B, <sup>s</sup>AUBERT 05A, <sup>t</sup>AITALA 97C, <sup>u</sup>LINK 04F, <sup>v</sup>AITALA 01A.

and  $\sigma_2$ . Second, following the methodology of FOCUS [LINK 04], by applying a  $K$ -matrix model to the  $\pi\pi$  S-wave [12].

Charm Dalitz-plot analyses might also prove useful for calibrating tools used to study  $B$  decays: specifically, to extract  $\alpha$  from  $B^0 \rightarrow \pi^+\pi^-\pi^0$ ,  $\beta$  from  $b \rightarrow s$  penguin decays (*e.g.*,  $B^0 \rightarrow \bar{K}^0 K^+ K^-$ ), and  $\gamma$  from  $B^\pm \rightarrow DK^\pm$  followed by  $D^0 \rightarrow \pi^+\pi^-\pi^0$  or  $K_S^0 K^+ K^-$  or  $K^+ K^- \pi^0$ , in addition to the well-studied  $D^0 \rightarrow K_S^0 \pi^+ \pi^-$  [2, 3].

**$D \rightarrow \pi\pi\pi$ :  $\sigma(500)$  or  $f_0(600)$ :** The decay  $D^+ \rightarrow \pi^+ \pi^+ \pi^-$  has been studied by the E687, E791 and FOCUS experiments (see Table 1). The E687 analysis considered the modes  $\rho(770)^0 \pi^+$ ,  $f_0(980) \pi^+$ ,  $f_2(1270) \pi^+$ , and a nonresonant component. E791 included, in addition,  $f_0(1370) \pi^+$  and  $\rho(1450)^0 \pi^+$ . Both analyses found a very large fraction ( $\sim 50\%$ ) for the nonresonant component, perhaps indicating a broad scalar contribution. E791 found the nonresonant amplitude to be consistent with zero if a broad scalar resonance was included. FOCUS analyzed its data using both the Breit-Wigner formalism and the  $K$ -matrix formalism for the  $\pi^+ \pi^-$  S-wave, following a 5-pole, 5-resonance model of Anisovich and Sarantsev [16]. The Breit-Wigner analysis included  $\rho(770)^0$ ,  $f_0(980)$ ,  $f_2(1270)^0$ ,  $f_0(1500)$ ,  $\sigma(500)$ , and a nonresonant component. The  $K$ -matrix formalism, with Breit-Wigner forms for the  $\rho(770)$  and  $f_2(1270)$ , also describe the FOCUS data well. None of these analyses has modeled the dynamics of the  $\pi^+ \pi^+$  interaction. Consideration of the  $I = 2$  S- and D-wave phase shifts, also measured in  $\pi^+ p \rightarrow \pi^+ \pi^+ n$  [18], could affect the  $\pi^+ \pi^-$  S-wave result.

Using the E791 data, Bediaga and Miranda [19] found additional evidence that the low-mass  $\pi^+ \pi^-$  feature is resonant by examining the phase of the  $\pi^+ \pi^-$  amplitude in the vicinity of the reported  $\sigma(500)$  mass. The phase variation with invariant mass is consistent with a resonant interpretation.

Table 1 gives the parameters of the  $\sigma(500)$  determined in charm Dalitz-plot analyses. A consistent relative phase between the  $\sigma(500)$  and  $\rho(770)$  resonances is observed.

**Table 2:** Parameters of the  $\sigma(500)$  resonance.  
The amplitude and phase are relative to the  $\rho(770)$ .

Experiment	E791 <sup>a</sup>	CLEO <sup>b</sup>	FOCUS [5]
Decay mode	$D^+ \rightarrow \pi^+\pi^+\pi^-$	$D^0 \rightarrow K_S^0\pi^+\pi^-$	$D^+ \rightarrow \pi^+\pi^+\pi^-$
Amplitude	$1.17 \pm 0.13 \pm 0.06$	$0.57 \pm 0.13$	—
Phase (°)	$205.7 \pm 8.0 \pm 5.2$	$214 \pm 11$	$200 \pm 31$
$m$ (MeV/c <sup>2</sup> )	$478_{-23}^{+24} \pm 17$	$513 \pm 32$	$443 \pm 27$
$\Gamma$ (MeV/c <sup>2</sup> )	$324_{-40}^{+42} \pm 21$	$335 \pm 67$	$443 \pm 80$

See the end of the  $D^+$  and  $D^0$  listings for these references:  
<sup>a</sup>AITALA 01B, <sup>b</sup>MURAMATSU 02.

CLEO has studied  $D^0 \rightarrow \pi^+\pi^-\pi^0$  (see Table 1). Only the three  $\rho(770)\pi$  resonant contributions are observed. No evidence is found for any  $\pi\pi$   $S$ -wave, either with the Breit-Wigner or with a  $K$ -matrix parametrization, using the 4-pole, 2-resonance model of Au, Morgan, and Pennington [17].

**$D^+ \rightarrow K^-\pi^+\pi^+$ : a  $\kappa(800)$ ?** Evidence for a broad  $K\pi$  scalar resonance has been found by E791 in  $D^+ \rightarrow K^-\pi^+\pi^+$  (see Table 1). Fitting the Dalitz plot with  $\bar{K}^*(892)^0\pi^+$ ,  $\bar{K}_0^*(1430)^0\pi^+$ ,  $\bar{K}_2^*(1430)^0\pi^+$ , and  $\bar{K}^*(1680)^0\pi^+$ , plus a constant nonresonant component, E791 found results consistent with earlier analyses by E691 and E687, with a nonresonant fit fraction of over 90%. With more events than the other experiments, E791 was then led to include an extra low-mass  $S$ -wave  $\bar{K}\pi$  resonance to account for the poor fit already seen by earlier experiments. A  $\kappa(800)$  with mass  $797 \pm 19 \pm 43$  MeV and width  $410 \pm 43 \pm 87$  MeV much improved the fit. The  $\kappa(800)$  became the dominant resonance and the nonresonant fit fraction was reduced from  $90.9 \pm 2.6\%$  to  $13.0 \pm 5.8 \pm 4.4\%$ .

In addition, E791 modeled the  $K\pi$   $S$ -wave phase variation as a function of  $K\pi$  mass with only the  $K_0^*(1430)$  resonance and a nonresonant component following a parametrization of LASS [20]. It was necessary to relax the unitarity constraint to describe the data [21]. The  $K\pi$   $S$ -wave phase behavior in this model was consistent with the model that included the  $\kappa$  resonance.

Finally, E791 performed a model-independent partial-wave analysis [AITALA 05] of the  $S$ -wave component of the  $K\pi$  system, finding the amplitude and phase from the  $K\pi$  threshold up to 1.72 GeV. No assumptions were made regarding dependence on invariant mass, but the analysis did use the relatively well-understood  $P$ - and  $D$ -waves, described by the  $K^*(892)$  and  $K^*(1680)$  and by the  $K_2(1430)$ , respectively. The results were similar to those obtained by AITALA 02, which parametrized the  $S$ -wave with  $\kappa$  and  $K_0(1430)$  Breit-Wigner forms and a constant complex non-resonant term. As with the  $\sigma(500)$ , the  $K^-\pi^+$   $S$ -wave result could be affected by including dynamics of the  $I = 2$   $\pi^+\pi^+$  interaction; however in AITALA 05, the  $I = 2$  elastic amplitude was found to be negligible compared to the  $\kappa$ .

CLEO allowed scalar  $K\pi$  resonances in fits to  $D^0 \rightarrow K^-\pi^+\pi^0$  and  $D^0 \rightarrow K_S^0\pi^+\pi^-$  (see Table 1), and observed a significant contribution from only the  $K_0^*(1430)$  [22]. BABAR fit  $D^0 \rightarrow K^0K^-\pi^+$  with both positively charged and neutral  $\bar{K}^*(892)$ ,  $\bar{K}_0^*(1430)$ ,  $\bar{K}_2^*(1430)$ , and  $\bar{K}^*(1680)$  resonances, as well as the  $a_0(980)^-$ ,  $a_0(1450)^-$ , and  $a_2(1310)^-$  resonances, and a nonresonant component [14]. BABAR also fit  $D^0 \rightarrow \bar{K}^0K^+\pi^-$  with the same resonances except for the  $a_2(1310)^-$ . In both cases, a good fit was obtained without including the  $\kappa$ .

FOCUS has conclusively observed a  $K\pi$   $S$ -wave as a distortion of the  $K^*(892)$  line-shape in semileptonic charm decays [LINK 02E, LINK 05D].

**The  $f_0(980)$ ,  $f_0(1370)$  and  $f_0(1500)$ :** The meson content of the  $0^{++}$  nonet and the quark content of the  $f_0(980)$ ,  $a_0(980)$ ,  $f_0(1370)$ ,  $f_0(1500)$ , and  $f_0(1710)$  mesons are current puzzles in light-meson spectroscopy [22]. Measuring branching fractions and couplings to different final states and comparing scalar-meson production rates among  $D^0$ ,  $D^+$ , and  $D_s^+$  mesons may help solve these puzzles.

For example: A large contribution of  $f_0(980)$  to  $D^0 \rightarrow K_S^0K^+K^-$  was reported by ARGUS [ALBRECHT 87E] and by BABAR [14]. This is inconsistent with the smaller contribution of  $f_0(980)$  observed in  $D^0 \rightarrow K_S^0\pi^+\pi^-$  by CLEO. The explanation is that  $D^0 \rightarrow K_S^0K^+K^-$  has a large contribution from  $a_0(980)^0 \rightarrow K^+K^-$ . Therefore CLEO studied

$D^0 \rightarrow K_S^0 \eta \pi^0$  [RUBIN 04], where the dominant contribution is from  $K_S^0 a_0(980)^0, a_0(980)^0 \rightarrow \eta \pi^0$ , and there can be no  $f_0(980)$ . A more recent BABAR analysis of  $D^0 \rightarrow K_S^0 K^+ K^-$  found a large amount of  $a_0(980) \rightarrow K \bar{K}$  and little  $f_0(980)$  [AUBERT 05B].

The proximity of the  $K \bar{K}$  threshold requires either a coupled-channel Breit-Wigner function [23] or a Flatté parametrization [24] of the  $f_0(980)$ . The width of the  $f_0(980)$  is poorly known. E791 and FOCUS [LINK 05C] [5] used a coupled-channel Breit-Wigner function to describe the  $f_0(980)$  in  $D_s^+ \rightarrow \pi^+ \pi^+ \pi^-$ . BESII studied the  $f_0(980)$  in  $J/\psi \rightarrow \phi \pi^+ \pi^-$  and  $\phi K^+ K^-$  [25]. The values found for the couplings to the  $\pi\pi$  and  $K \bar{K}$  channels,  $g_{\pi\pi}$  and  $g_{KK}$ , were not consistent. Results such as these are desirable for input to the analysis of  $D_s^+ \rightarrow K^+ K^- \pi^+$  [15], which includes both the  $f_0(980)$  and  $a_0(980)$ .

The quark content of the  $f_0(1370)$  and  $f_0(1500)$  can perhaps be inferred from how they populate various Dalitz plots. Results so far are confusing. The E791 analysis of  $D^+ \rightarrow \pi^+ \pi^+ \pi^-$  [AITALA 01B] found some  $f_0(1370)$  but no  $f_0(1500)$ , while the FOCUS analysis [5] of this mode found little  $f_0(1370)$ . In  $D_s^+ \rightarrow \pi^+ \pi^+ \pi^-$ , E687 and FOCUS [5] found no  $f_0(1370)$ , but did find a resonance with parameters similar to the  $f_0(1500)$ , whereas E791 found a  $\pi^+ \pi^-$  resonance with mass  $1434 \pm 18 \pm 9$  MeV and width  $172 \pm 32 \pm 6$  MeV, consistent with neither the  $f_0(1370)$  or  $f_0(1500)$ . BABAR [AUBERT 05B] in  $D^0 \rightarrow \bar{K}^0 K^+ K^-$  found neither the  $f_0(1370)$  nor the  $f_0(1500)$ , but did observe a  $K^+ K^-$  resonance consistent with the values from E791 given above, while CLEO has observed the  $f_0(1370)$  in  $D^0 \rightarrow K_S^0 \pi^+ \pi^-$ . The FOCUS analysis that used the  $K$ -matrix formalism for the  $\pi\pi$  S-wave observed significant couplings to five  $T$ -matrix poles— $f_0(980), f_0(1300), f_0(1200-1600), f_0(1500), f_0(1750)$ —in both  $D^+ \rightarrow \pi^+ \pi^- \pi^+$  and  $D_s^+ \rightarrow \pi^+ \pi^- \pi^+$ . Again, the quark content of each pole might be inferred from the coupling to various Dalitz plots.

It is noteworthy that the S-wave observed in  $B$  Dalitz-plot analyses appears to be different than that observed in  $D$ -meson decays.

**Doubly Cabibbo-Suppressed Decays:** There are two classes of multibody doubly Cabibbo-suppressed (DCS) decays of  $D$  mesons. The first consists of those in which the DCS and corresponding Cabibbo-favored (CF) decays populate distinct Dalitz plots: the pairs  $D^0 \rightarrow K^+\pi^-\pi^0$  and  $D^0 \rightarrow K^-\pi^+\pi^0$ , or  $D^+ \rightarrow K^+\pi^+\pi^-$  and  $D^+ \rightarrow K^-\pi^+\pi^+$ , are examples. CLEO [BRANDENBURG 01] and Belle [TIAN 05] have reported  $\Gamma(D^0 \rightarrow K^+\pi^-\pi^0)/\Gamma(D^0 \rightarrow K^-\pi^+\pi^0) = (0.43_{-0.10}^{+0.11} \pm 0.07)\%$  and  $(0.229 \pm 0.015_{-0.009}^{+0.013})\%$ , respectively. E791 and FOCUS have reported  $\Gamma(D^+ \rightarrow K^+\pi^-\pi^+)/\Gamma(D^+ \rightarrow K^-\pi^+\pi^+) = (0.77 \pm 0.17 \pm 0.08)\%$  and  $(0.65 \pm 0.08 \pm 0.04)\%$ , respectively.

The second class consists of decays in which the DCS and CF modes populate the same Dalitz plot; for example,  $D^0 \rightarrow K^{*-}\pi^+$  and  $D^0 \rightarrow K^{*+}\pi^-$  both contribute to  $D^0 \rightarrow K_S^0\pi^+\pi^-$ . In this case, the potential for interference of DCS and CF amplitudes increases the sensitivity to the DCS amplitude. CLEO found the relative amplitude and phase to be  $(7.1 \pm 1.3_{-0.6}^{+2.6} {}_{-0.6}^{+2.6})\%$  and  $(189 \pm 10 \pm 3_{-5}^{+15})^\circ$ , corresponding to  $\Gamma(D^0 \rightarrow K^*(892)^+\pi^-)/\Gamma(D^0 \rightarrow K^*(892)^-\pi^+) = (0.5 \pm 0.2_{-0.1}^{+0.5} {}_{-0.1}^{+0.4})\%$ . In addition to  $D^0 \rightarrow K^*(892)^+\pi^-$ , Belle [10,11] and

BABAR [12,13] have found evidence for  $D^0 \rightarrow K_0(1430)^+\pi^-$  and  $K_2(1430)^+\pi^-$ , and Belle has also found evidence for  $K^*(1680)^+\pi^-$ .

**$CP$  Violation:** In the limit of  $CP$  conservation, charge conjugate decays will have the same Dalitz-plot distribution. The  $D^{*\pm}$  tag enables the discrimination between  $D^0$  and  $\overline{D}^0$ . The integrated  $CP$  violation across the Dalitz plot is determined in two ways. The first uses

$$\mathcal{A}_{CP} = \int \left( \frac{|\mathcal{M}|^2 - |\overline{\mathcal{M}}|^2}{|\mathcal{M}|^2 + |\overline{\mathcal{M}}|^2} \right) dm_{ab}^2 dm_{bc}^2 \bigg/ \int dm_{ab}^2 dm_{bc}^2, \quad (1)$$

where  $\mathcal{M}$  and  $\overline{\mathcal{M}}$  are the  $D^0$  and  $\overline{D}^0$  Dalitz-plot amplitudes. The second uses the asymmetry in the efficiency-corrected  $D^0$  and  $\overline{D}^0$  yields,

$$\mathcal{A}_{CP} = \frac{N_{D^0} - N_{\overline{D}^0}}{N_{D^0} + N_{\overline{D}^0}}. \quad (2)$$

These expressions are less sensitive to  $CP$  violation than are the individual resonant submodes [ASNER 04A]. Table 3 lists the results for  $CP$  violation. No evidence of  $CP$  violation has been observed in  $D$ -meson decays.

**Table 3:** Dalitz-plot-integrated  $CP$  violation. Measurements computing  $\mathcal{A}_{CP}$  with Eq. (2) rather than Eq. (1) are denoted  $^\dagger$ .

Experiment	Decay mode	$\mathcal{A}_{CP}(\%)$
BABAR <sup>a</sup>	$D^+ \rightarrow K^+ K^- \pi^+$	$1.4 \pm 1.0 \pm 0.8$
Belle <sup>b\dagger</sup>	$D^0 \rightarrow K^+ \pi^- \pi^0$	$-0.6 \pm 5.3$
Belle <sup>b\dagger</sup>	$D^0 \rightarrow K^+ \pi^- \pi^+ \pi^-$	$-1.8 \pm 4.4$
CLEO <sup>c</sup>	$D^0 \rightarrow K^- \pi^+ \pi^0$	$-3.1 \pm 8.6$
CLEO <sup>d\dagger</sup>	$D^0 \rightarrow K^+ \pi^- \pi^0$	$+9^{+22}_{-25}$
CLEO <sup>e</sup>	$D^0 \rightarrow K_S^0 \pi^+ \pi^-$	$-0.9 \pm 2.1^{+1.0+1.3}_{-4.3-3.7}$
CLEO <sup>f</sup>	$D^0 \rightarrow \pi^+ \pi^- \pi^0$	$+1^{+9}_{-7} \pm 9$

See the end of the  $D^+$  and  $D^0$  Listings for these references:

<sup>a</sup>AUBERT 05A, <sup>b</sup>TIAN 05, <sup>c</sup>KOPP 01, <sup>d</sup>BRANDENBURG 01, <sup>e</sup>ASNER 04A, <sup>f</sup>CRONIN-HENNESSY 05.

The possibility of interference between  $CP$ -conserving and  $CP$ -violating amplitudes provides a more sensitive probe of  $CP$  violation. The constraints on the square of the  $CP$ -violating amplitude obtained in the resonant submodes of  $D^0 \rightarrow K_S^0 \pi^+ \pi^-$  range from  $3.5 \times 10^{-4}$  to  $28.4 \times 10^{-4}$  at 95% confidence level [ASNER 04A].

## References

1. See the note on “ $D^0$ - $\overline{D}^0$  Mixing” in this *Review*.
2. See the note on “The CKM Quark Mixing Matrix” in this *Review*.
3. See the note on “ $CP$  Violation in Meson Decays” in this *Review*.



4. Dalitz plot analysis of the wrong sign rate  $D^0 \rightarrow K^+\pi^-\pi^0$  [BRANDENBURG 01] and the time dependence of Dalitz plot analysis of  $D^0 \rightarrow K_S^0\pi^+\pi^-$  [ASNER 05] are two candidate processes.
5. S. Malvezzi, AIP Conf. Proc. **688**, 276 (2004) [Nucl. Phys. Proc. Suppl. **126**, 220 (2004)].
6. A. Giri *et al.*, Phys. Rev. **D68**, 054018 (2003).
7. A. Bondar and A. Poluektov, hep-ph/0510246.
8. J. Blatt and V. Weisskopf, *Theoretical Nuclear Physics*, New York: John Wiley & Sons (1952).
9. F. von Hippel and C. Quigg, Phys. Rev. **D5**, 624, (1972).
10. A. Poluektov *et al.* (Belle Collab.), Phys. Rev. D **70**, 072003 (2004).
11. K. Abe *et al.* Belle Collaboration, hep-ex/0411049.
12. B. Aubert *et al.* (BABAR Collab.), Phys. Rev. Lett. **95**, 121802 (2005).
13. B. Aubert *et al.* (BABAR Collab.), hep-ex/0507101.
14. B. Aubert *et al.* (BABAR Collab.), hep-ex/0207089; contributed to the 31st International Conference on High Energy Physics (ICHEP 2002).
15. S. Malvezzi, AIP Conf. Proc. **549**, 569 (2002).
16. V. V. Anisovich and A. V. Sarantsev, Eur. Phys. J. A **16**, 229 (2003).
17. K. L. Au *et al.*, Phys. Rev. D **35**, 1633 (1987).
18. W. Hoogland *et al.*, Nucl. Phys. **B69**, 266 (1974).
19. I. Bediaga (E791 Collab.), AIP Conf. Proc. **688**, 252 (2004).
20. D. Aston *et al.* (LASS Collab.), Nucl. Phys. **B296**, 493 (1988).
21. C. Gobel (Fermilab E791 Collab.), AIP Conf. Proc. **688**, 266 (2004).
22. See the “Note on Scalar Mesons” in this *Review*.
23. T. A. Armstrong *et al.* (WA76 Collab.), Z. Phys. **C51**, 351 (1991).
24. S. M. Flatte, Phys. Lett. **B63**, 224 (1976).
25. M. Ablikim *et al.* (BES Collab.), Phys. Lett. B **607**, 243 (2005).